

ODE Refresher

1st-order linear

$$y'(t) + p(t)y(t) = g(t)$$

integrating factor

$$\mu(t) = e^{\int p(t) dt}$$

leading coefficient is 1

multiply the ODE by $\mu(t)$

$$\mu y' + p\mu y = g\mu$$

$\underbrace{\hspace{10em}}$

$$\frac{d}{dt}(\mu y) = \mu g$$

integrate both sides and solve for y

example

$$y' + \frac{2}{t}y = \frac{\cos(t)}{t^2} \quad (t > 0)$$

$$\mu = e^{\int \frac{2}{t} dt} = e^{2 \ln t} = e^{\ln t^2} = t^2$$

$$t^2 \left(y' + \frac{2}{t} y \right) = t^2 \cdot \frac{\cos(t)}{t^2}$$

$$t^2 y' + 2t y = \cos(t)$$



$$\frac{d}{dt} (t^2 y) = \cos(t)$$

integrate both sides

$$t^2 y = \sin(t) + C$$

$$y = \frac{\sin(t)}{t^2} + \frac{C}{t^2}$$

it will show up later when we solve the heat equation

$$\frac{\partial u}{\partial t} = \alpha^2 \frac{\partial^2 u}{\partial x^2} \quad u = u(x, t)$$

there is a 1st-order linear that needs to be solved

$$T'(t) + \alpha^2 \lambda T(t) = 0$$

α, λ constants > 0

this can also be solved as a separable eq.

$$\frac{dT}{dt} = -\alpha^2 \lambda T$$

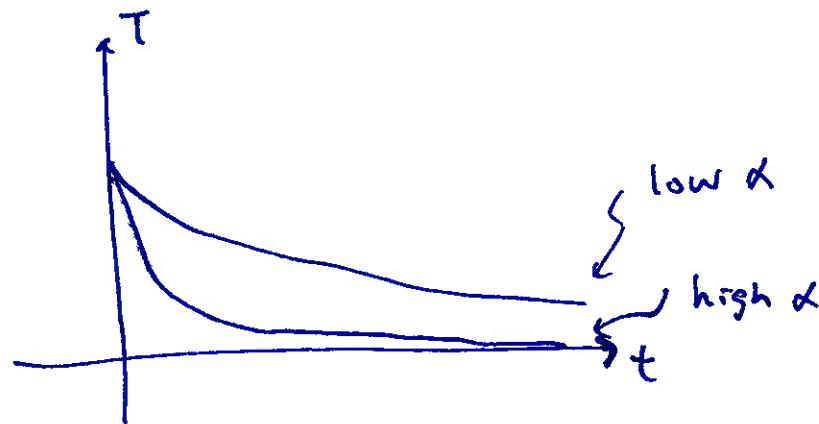
$$\frac{1}{T} dT = -\alpha^2 \lambda dt$$

$$\ln|T| = -\alpha^2 \lambda t + C$$

$$T(t) = e^{-\alpha^2 \lambda t + C} = e^{-\alpha^2 \lambda t} (e^C) \quad \text{constant}$$

$$T(t) = C e^{-\alpha^2 \lambda t}$$

in the context of heat eq, this tells us how heat dissipates as a function of time for a given λ and position.



2nd-order homogeneous

$$y'' + ky = 0$$

Solutions are of the form e^{rt}

r is solution of the characteristic eq.

$$r^2 + k = 0$$

$$r^2 = -k$$

if $k = 0$, $r = 0, 0$, $y = C_1 + C_2 t$ (line)

if $k > 0$, $r = \pm \sqrt{k} = \pm i\sqrt{|k|}$ $y = \cos(\sqrt{k}t)$, $y = \sin(\sqrt{k}t)$

$$y = C_1 \cos(\sqrt{k}t) + C_2 \sin(\sqrt{k}t) \quad (\text{bounded})$$

if $k < 0$, $r = \pm \sqrt{|k|}$ $y = e^{\sqrt{|k|}t}$, $y = e^{-\sqrt{|k|}t}$

$$y = C_1 e^{\sqrt{|k|}t} + C_2 e^{-\sqrt{|k|}t} \quad (\text{unbounded})$$

$$y = d_1 \cosh(\sqrt{k}t) + d_2 \sinh(\sqrt{k}t)$$

$$\cosh(x) = \frac{e^x + e^{-x}}{2} \quad \sinh(x) = \frac{e^x - e^{-x}}{2}$$

this shows up when we solve the heat, wave, and Laplace's eqs.

$\xrightarrow{\text{uppercase } x}$
 $\bar{X}''(x) + \lambda \bar{X}(x) = 0$
 or $Y''(y) + \lambda Y(y) = 0$

solutions: $\bar{X}(x) = C_1 \cos(\sqrt{\lambda} x) + C_2 \sin(\sqrt{\lambda} x)$ if $\lambda \geq 0$

$\bar{X}(x) = C_1 \cosh(\sqrt{\lambda} x) + C_2 \sinh(\sqrt{\lambda} x)$ if $\lambda < 0$

$\bar{X}(x) = C_1 + C_2 x$ if $\lambda = 0$

in the context of heat eq, this governs the space part of the solution (how temperature varies w/ position if time is fixed)

now briefly about nonhomogeneous 2nd-order

in the context of mass-spring system: $m x'' + k x = f(t)$

when $f(t) = F_0 \sin(\omega t)$ (sinusoidal input)

ω : input frequency

F_0 : input magnitude

w/o input: $m x'' + k x = 0$

$x(t) = C_1 \cos\left(\sqrt{\frac{k}{m}} t\right) + C_2 \sin\left(\sqrt{\frac{k}{m}} t\right)$ (complementary solution)

w/ input: $x(t) = C_1 \cos\left(\sqrt{\frac{k}{m}} t\right) + C_2 \sin\left(\sqrt{\frac{k}{m}} t\right) + x_p$

↑
particular solution

$\sqrt{\frac{k}{m}} = \omega_0$ (natural frequency)

if $\omega \neq \omega_0$ (input freq \neq natural freq)

we can solve using method of undetermined coefficients

$$mX'' + kX = F_0 \sin(\omega t)$$

$$X_p = A \cos(\omega t) + B \sin(\omega t)$$

Sub into the ODE

∴

$$A = 0, \quad B = \frac{F_0}{k - m\omega^2} \quad \text{trouble if } k = m\omega^2 \text{ or } \omega = \sqrt{\frac{k}{m}}$$

$$X(t) = C_1 \cos\left(\sqrt{\frac{k}{m}} t\right) + C_2 \sin\left(\sqrt{\frac{k}{m}} t\right) + \frac{F_0}{k - m\omega^2} \sin(\omega t)$$

but if input freq = natural freq ($\omega = \sqrt{\frac{k}{m}}$, $k - m\omega^2 = 0$)

X_p needs to be adjusted due to duplication of complementary

$$X_p = A \underline{t} \cos(\omega t) + B \underline{t} \sin(\omega t)$$

Sub into $mX'' + kX = F_0 \sin(\omega t)$

∴

$$B = 0, \quad A = -\frac{F_0}{2m\omega}$$

$$x(t) = C_1 \cos\left(\sqrt{\frac{k}{m}} t\right) + C_2 \sin\left(\sqrt{\frac{k}{m}} t\right) - \frac{F_0}{2m\omega} t \cos(\omega t)$$

↑ goes to ∞ as t grows

resonance